The radiation exposure to patients undergoing CT examinations is determined by two factors: equipment-related factors, i.e. the design of the scanner with respect to dose efficiency, and application-related factors, i.e. the way in which the radiologist or the radiographer makes use of the scanner. In this chapter, the features and parameters influencing patient dose are outlined. First, however, a brief introduction on the dose descriptors applicable to CT is given.

### 3.1 CT Dose Descriptors

The dose quantities used in projection radiography are not applicable to CT for three reasons:

- **First**, the dose distribution inside the patient is completely different from that for a conventional radiogram, where the dose decreases continuously from the entrance of the X-ray beam to its exit, with a ratio of between 100 and 1000 to 1. In the case of CT, as a consequence of the scanning procedure that equally irradiates the patient from all directions, the dose is almost equally distributed in the scanning plane. A dose comparison of CT with conventional projection radiography in terms of skin dose therefore doesn’t make any sense.

- **Second**, the scanning procedure using narrow beams along the longitudinal \( z \)-axis of the patient implies that a significant portion of the radiation energy is deposited outside the nominal beam width. This is mainly due to penumbra effects and scattered radiation produced inside the beam.

- **Third**, the situation in CT is further complicated by the circumstances in which - unlike in conventional projection radiography - the volume to be imaged is not irradiated simultaneously. This often leads to confusion about what the dose from a complete series of e.g. 15 slices might be compared with the dose from a single slice.

As a consequence, dedicated dose quantities that account for these peculiarities are needed: The ‘Computed Tomography Dose Index (CTDI)’, which is a measure of the local dose, and the ‘Dose-Length Product (DLP)’, representing the integral radiation exposure associated with a CT examination. Fortunately, a bridge exists that enables to compare CT with radiation exposure from other modalities and sources; this can be achieved by the effective dose (E). So there are three dose descriptors in all, which everyone dealing with CT should be familiar with.

#### 3.1.1 Computed Tomography Dose Index

The ‘Computed Tomography Dose Index (CTDI)’ is the fundamental CT dose descriptor. By making use of this quantity, the first two peculiarities of CT scanning are taken into account: The CTDI (unit: Milligray (mGy)) is derived from the dose distribution along a line which is parallel to the axis of rotation for the scanner (= \( z \)-axis) and which is recorded for a single rotation of the x-ray source. Fig. 3.1 illustrates the meaning of this term: CTDI is the equivalent of the dose value inside the irradiated slice (beam) that would result if the absorbed radiation dose profile were entirely concentrated to a rectangular profile of width equal to the nominal beam width \( N \cdot h_{col} \), with \( N \) being the number of independent (i.e. non-overlapping) slices that are acquired simultaneously. Accordingly, all dose contributions from outside the nominal beam width, i.e. the areas under the tails of the dose profile, are added to the area inside the slice.

**Fig. 3.1 Illustration of the term ‘Computed Tomography Dose Index (CTDI)’**:
CTDI is the equivalent of the dose value inside the irradiated slice (beam) that would result if the absorbed radiation dose profile were entirely concentrated to a rectangular profile of width equal to the nominal beam width \( N \cdot h_{col} \).
Chapter 3: CT Parameters that Influence the Radiation Dose

The corresponding mathematical definition of CTDI therefore describes the summation of all dose contributions along the z-axis:

\[
CTDI = \frac{1}{N \cdot h_{\text{col}}} \cdot \int_{-\infty}^{+\infty} D(z) \cdot dz
\]  

(3.1)

where \( D(z) \) is the value of the dose at a given location, \( z \), and \( N \cdot h_{\text{col}} \) is the nominal value of the total collimation (beam width) that is used for data acquisition. CTDI is therefore equal to the area of the dose profile (the ‘dose-profile integral’) divided by the nominal beam width. In practice, the dose profile is accumulated in a range of –50 mm to +50 mm relative to the centre of the beam, i.e. over a distance of 100 mm.

The relevancy of CTDI becomes obvious from the total dose profile of a scan series with e.g. \( n=15 \) subsequent rotations (fig. 3.2). The average level of the total dose profile, which is called ‘Multiple Scan Average Dose (MSAD)’ (Shope 1981), is higher than the peak value of each single dose profile. This increase results from the tails of the single dose profiles for a scan series. Obviously, MSAD and CTDI are exactly equal if the table feed TF is equal to the nominal beam width \( N \cdot h_{\text{col}} \), i.e. if the pitch factor

\[
p = \frac{TF}{N \cdot h_{\text{col}}}
\]  

(3.2)

is equal to 1. In general (i.e. if the pitch is not equal to 1, see fig. 3.3), the relationship between CTDI and MSAD is given by

\[
MSAD = \frac{1}{p} \cdot CTDI
\]  

(3.3)

The practical implication of equation (3.3) is that - in order to obtain the average dose for a scan series - it is not necessary to carry out all the scans. Instead, it is sufficient to obtain the CTDI from a single scan by acquiring the entire dose profile according to equation (3.1). This is achieved with dose measurements using long, pencil-like detectors, with an active length of 10 cm (fig. 3.4). These detectors accumulate the dose profile integral (DPI, unit: mGy·cm), i.e. the area under the dose profile shown in fig. 3.1. The CTDI is then obtained according to equation 3.1 by division with the nominal beam width \( N \cdot h_{\text{col}} \).

In order to obtain estimates of the dose to organs that are located in the scan range, the CTDI generally refers to standard dosimetry phantoms with patient-like diameters. In the standard measuring procedure for CTDI, which utilizes two cylindrical Perspex (PMMA) phantoms of different diameter (fig. 3.4), dose is measured at the centre and near the periphery of the phantom (fig. 3.5). The larger phantom, being 32 cm in diameter, represents the absorption that is typical for the trunk region of adults. The smaller phantom (16 cm in diameter) represents the patient in head examinations. The smaller phantom is also used for dose assessment in paediatric examinations (Shrimpton 2000). The dose values thus obtained are denoted as

\[
CTDI_{\text{H,e}} \quad \text{and} \quad CTDI_{\text{H,p}}
\]  

(3.4a)

\[\text{Fig. 3.2 Total dose profile of a scan series with } n=15 \text{ subsequent rotations. The average level of the total dose profile, which is called ‘Multiple Scan Average Dose (MSAD)’, is equal to CTDI if the table feed TF is equal to the nominal beam width } N \cdot h_{\text{col}} \text{ (i.e. pitch } p = 1).\]

\[\text{Fig. 3.3 Total dose profile of a scan series with } n=15 \text{ subsequent rotations, scanned with pitch } = 0.7, \text{ however. Due to the larger overlap, MSAD is higher than in fig. 3.2 and amounts to CTDI divided by pitch.}\]
with $H =$ head, $B =$ body, $c =$ centre, $p =$ periphery.

To make life easier, each pair of CTDI values (central and peripheral) can be combined into one single one named ‘Weighted CTDI (CTDI<sub>w</sub>)’, which represents the CTDI averaged over the cross section of the pertaining phantom: 

$$CTDI_{w} = \frac{1}{3} \cdot CTDI_{c} + \frac{2}{3} \cdot CTDI_{p}$$  \hspace{1cm} (3.5)

where the subscript XYZ stands either for H(ead) or B(ody). In daily practice, CTDI<sub>w</sub> is used as one of two dose descriptors for dose recommendations (‘reference values’) that have been introduced by the European Commission (1999a).

If pitch-related effects on the radiation exposure are taken into account at the level of local dose (i.e. CTDI) already, a quantity named ‘Volume CTDI (CTDI<sub>vol</sub>)’ is defined (IEC 2001):

$$CTDI_{vol} = \frac{CTDI}{p}$$  \hspace{1cm} (3.6)

So CTDI<sub>vol</sub> is the pitch-corrected CTDI<sub>w</sub>. Apart from the integration length, which is limited to 100 mm, CTDI<sub>vol</sub> is practically identical to MSAD based on CTDI<sub>w</sub> (i.e. MSAD<sub>w</sub>). Since averaging includes both the cross section and the scan length, CTDI<sub>vol</sub> therefore represents the average dose for a given scan volume. CTDI<sub>vol</sub> is used as the dose quantity that is displayed at the operator’s console of newer scanners. This also holds true even if the display is labelled as ‘CTDI<sub>w</sub>’ due to faulty definition in the first edition of the particular IEC standard for CT (IEC 1999), or simply as ‘CTDI’.

Attention is required if the dose displayed as CTDI<sub>vol</sub> shall be used for comparison with reference values given in terms of CTDI<sub>w</sub>. For this purpose, the pitch correction introduced in equation (3.6) needs to be reversed by multiplying the CTDI<sub>vol</sub> value with the pitch factor. Care is also required if the CTDI<sub>vol</sub> displayed is used to assess paediatric radiation exposure: whether head or body CTDI values are displayed depends only on the scan mode (head or body), but not on the patient size. Consequently, the dose to children and infants undergoing CT examinations of the trunk region, which for the same scan parameter settings depends on the patient diameter, is currently underestimated with the dose displayed at the operator’s console by a factor 2 to 3.

CTDI statements in scanner specification sheets are given for the head phantom as well as for the body phantom, and often apply to a current-time product of 100 mAs or 1 mAs. In this case it must be recognized that a quantity named ‘normalized CTDI’ is used, which is labelled ‘CTDI (unit: mGy/mAs)’ in order to avoid confusion. The normalized CTDI is obtained by dividing the CTDI value by the mAs product $Q$ that was used to measure CTDI:

$$nCTDI = \frac{CTDI}{Q}$$  \hspace{1cm} (3.7)

It is worthwhile (and indeed necessary) to note that the normalized CTDI is a characteristic quantity for a scanner (dose rate coefficient) which simply represents the capacity of a scanner in terms of output and which conveys absolutely nothing about patient dose. Very often it is assumed that scanners with a high value of $nCTDI$ are more ‘dangerous’ than other models with lower $nCTDI$ values. This is not necessarily the case. Reference to patient dose cannot be made unless the normalized CTDI has been multiplied by the tube current-time product $Q$ that is required in order to produce images of diagnostic quality with the type of
scanner under consideration. Only after having carried out this step is it possible to decide if a particular scanner needs more or less dose than another model for a specified type of examination.

### 3.1.2 Dose-Length Product

The third peculiarity of CT, i.e. the question what the dose from a complete series of e.g. 15 slices might be compared with the dose from a single slice, is solved by introducing a dose descriptor named ‘dose-length product (DLP; unit: mGy·cm)’. DLP takes both the ‘intensity’ (represented by the CTDI\text{vol}) and the extension (represented by the scan length L) of an irradiation into account (fig. 3.6):

\[
DLP = CTDI_{\text{vol}} \cdot L
\]  

(3.8)

So the dose-length product increases with the number of slices (correctly: with the length of the irradiated body section), while the dose (i.e. CTDI\text{vol}) remains the same regardless of the number of slices or length, respectively. In fig. 3.6, the area of the total dose profile of the scan series represents the DLP. DLP is the equivalent of the dose-area product (DAP) in projection radiography, a quantity that also combines both aspects (intensity and extension) of patient exposure.

In sequential scanning, the scan length is determined by the beam width \(N \cdot h_{\text{col}}\) and the number \(n\) of tables feeds TF:

\[
L = n \cdot TF + N \cdot h_{\text{col}}
\]  

(3.9)

while in spiral scanning the scan length only depends on the number \(n\) of rotations and the table feed TF:

\[
L = n \cdot TF = \frac{T}{t_{\text{rot}}} \cdot p \cdot N \cdot h_{\text{col}}
\]  

(3.10)

where \(T\) is the total scan time, \(t_{\text{rot}}\) is the rotation time, and \(p\) is the pitch factor. While in sequential scanning the scan length \(L\) is equal to the range from the begin of the first slice till the end of the last, the (gross) scan length for spiral scanning not only comprises the (net) length of the imaged body section but also includes the additional rotations at the begin and the end of the scan (‘over-ranging’) that are required for data interpolation.

If an examination consists of several sequential scan series or spiral scans, the dose-length product of the complete examination (DLP\text{exam}) is the sum of the dose-length products of each single series or spiral scan:

\[
DLP_{\text{exam}} = \sum_i DLP_i
\]  

(3.11)

In daily practice, the DLP is used as the second (and most important) of the two dose descriptors for dose recommendations (‘reference values’) that have been introduced by the European Commission (1999a).

### 3.1.3 Effective Dose

CTDI and DLP are CT-specific dose descriptors that do not allow for comparisons with radiation exposures from other sources, e.g. projection radiography, nuclear medicine or natural background radiation. The only common denominator to achieve this goal is the ‘Effective Dose’. With effective dose, the organ doses from a partial irradiation of the body are converted into an equivalent uniform dose to the entire body.

Effective dose \(E\) (unit: Millisievert (mSv)) according to ICRP 60 (ICRP 1991) is defined as the weighted average of organ dose values \(H_i\) for a number of specified organs:

\[
E = \sum_i w_i \cdot H_{T,i}
\]  

(3.12)

How much a particular organ contributes to effective dose depends on its relative sensitivity for radiation-induced effects, as represented by the tissue-weighting factor \(w_i\) attributed to the organ:

- 0.20 for gonads;
- 0.12 for each of lungs, colon, red bone marrow and stomach wall;
- 0.05 for each of breast, urinary bladder, liver, thyroid and oesophagus;
- 0.01 for each of skeleton and skin;
- 0.05 for the ‘remainder’

The ‘remainder’ consists of a group of additional organs and tissues with a lower sensitivity for radiation induced
effects for which the average dose must be used: small intestine, brain, spleen, muscle tissue, adrenals, kidneys, pancreas, thymus and uterus. The sum of all tissue-weighting factors \( w_i \) is equal to 1.

Effective dose cannot as such be measured directly in vivo. Measurements in anthropomorphic phantoms with thermoluminescent dosemeters (TLD) are very time-consuming and therefore not well suited for daily practice. Effective dose, however, can be assessed in various ways by using conversion factors. For coarse estimates, it is sufficient to multiply the dose-length product with mean conversion factors, depending on which one out of three body regions was scanned and whether the scan was made in head or body scanning mode:

\[
E = DLP \cdot f_{\text{mean}} \quad (3.13)
\]

For adults of standard size, the following generic mean conversion factors \( f_{\text{mean}} \) apply:

- 0.025 mSv/mGy·cm for the head region
- 0.060 mSv/mGy·cm for the neck region, scanned in head mode
- 0.100 mSv/mGy·cm for the neck region, scanned in body mode
- 0.175 mSv/mGy·cm for the trunk region

Similar factors \( E_{\text{DLP}} \), which additionally distinguish between chest, abdomen and pelvis, but do not account for differences in scan mode, are given in report EUR 16262 (European Commission 1999b).

In order to apply equation (3.13), the DLP or at least the CTDI \( \text{vol} \) and the (gross) scan length \( L \), from which the DLP can be calculated according to equation (3.8), must be available. If the scanner is not equipped with a dose display, or if a more detailed assessment of effective dose is desired (e.g. to be more specific for the scanned region of the body, to distinguish between males and females, to assess paediatric doses, or to take differences between scanners into account), dedicated CT dose calculation software should be used. These programs make use of more detailed conversion factors and also allow for calculation of organ doses. Currently, five different programs are either commercially available or in general use, which differ significantly in specifications, performance, and price. A comparative study on these programs has recently been published by Tack (2006).

Typical tolerances in effective dose assessment with these programs are in the order of ±20 to ±30%. Similar uncertainties also apply to effective dose assessment with TLD measurements in Alderson phantoms. This should always be borne in mind when comparing doses from different scanners in terms of effective dose. Care is also needed to not mix up effective dose with organ doses, as both are expressed in mSv. Nevertheless, effective dose is of great value, e.g. to answer questions raised from patients. For this purpose, the annual natural background radiation, which is between 2 and 3 mSv in many countries, can be used as a scale.

A comprehensive compilation of dose-relevant scanner data and other useful information required for CT dose assessment can be found in Nagel et al. (2002). The data given there apply to most of the scanners currently in use except for the most recent ones. However, data for these new scanners can be found in the CT-Expo software package (Stamm and Nagel 2001) that is based on the data and formalism outlined in this book and is updated regularly.

### 3.2 Equipment-related Factors

#### 3.2.1 Beam Filtration

In conventional projection radiography, beam filtration is a well-known means to reduce those portions of the radiation spectrum with no or little contribution to image formation. In the early years of CT history, beam filtration was comparatively large in order to compensate for beam-hardening artefacts. Filters made from 0.5 mm of copper, with filtering properties equivalent to approximately 18 mm of aluminium (quality equivalent filtration, Nagel 1986), were not unusual at that time. The present generation of scanners typically employs a beam filtration for the X-ray tube assembly of between 1 and 3 mm Al and an additional filtration (flat filter) of 0.1 mm Cu, giving a total beam filtration of between 5 to 6 mm Al.

Apart from this, there are a number of older and also newer scanners which operate with an added filtration of approximately 0.2 mm Cu, resulting in a total beam filtration of between 8 and 9 mm Al, and sometimes even more (currently up to 12 mm Al quality-equivalent filtration). Likewise, there are also scanners that employ less filtration. Consequently, the normalised dose values for these scanners (\( \text{CTDI}_{\text{vol}} \) in terms of mGy/mAs) differ significantly. Very often these lower or higher values are misunderstood as being an indicator that the equipment is more or less dose-efficient compared with other scanners. This might not necessarily be the case in reality.

Apart from dose, the consequences on image quality arising from the beam hardening and beam attenuating...
properties of filtration have also to be considered (Nagel 1989). The use of additional filtration impairs primary contrast and increases noise due to reduced beam intensity per mAs as experienced by the detectors. Without compensating for these adverse effects (e.g. by increasing tube current-time product), the contrast-to-noise ratio, which affects the detectability of small or low-contrast details, is reduced. Unpublished studies by the author show that, in order to maintain the contrast-to-noise ratio (i.e. for constant image quality), the net reduction in terms of effective dose achieved by increasing the standard beam filtration (1 mm Al + 0.1 mm Cu = 4.5 mm Al quality equivalent) by 0.2 mm Cu amounts to not more than 10%, even in favourable situations (soft tissue imaging, see fig. 3.7). Conversely, the same added filtration leads to higher patient doses (up to 15%) in examinations with administration of contrast agents (iodine). At the same time, tube loading must be increased by 20% in order to compensate for reduced beam intensity.

Newer surveys on CT practice (e.g. Galanski et al. 2001) revealed that scanners of comparable age but with largely differing beam filtration are operated at almost similar dose levels. Similar results in terms of dose efficiency have been found in comparative tests on scanners with differing beam filtration conducted by ImPACT (2004). Contrary to projection radiography, which operates at comparatively lower tube potentials, beam filtration plays only a minor role in CT where higher tube potentials are applied. A return to increased beam filtration - as sometimes recommended or practised - is less advantageous than expected and should only be made if sufficient X-ray tube loading capacity is available or if other important aspects exist (e.g. improved performance of reconstruction filters).

### 3.2.2 Beam Shaper

Most scanners are equipped with a dedicated filter device named ‘beam shaper’ or ‘bow-tie filter’ that modifies the spatial distribution of radiation emitted within the fan beam. The purpose of this kind of filter (which is characterised by increasing thickness towards its outer edges) is to adapt the beam intensity to match the reduced attenuation of objects in the outer portions of the fan beam.

---

**Fig. 3.8** Dose profiles free-in-air with umbra (dark grey) and penumbra (light grey) portions for a single-slice scanner (a.), a dual-slice scanner (b.), and a quad-slice scanner (c.). With single- and dual-slice scanners, the width of the active detector rows is sufficient to capture the entire dose profile, penumbra included (except for some scanners which employ restrictive post-patient collimation). For scanners with four and more slices acquired simultaneously, penumbra is excluded from detection in order to serve all detector channels equally well. The combined width of the penumbra triangles at both sides is characterized by the overbeaming parameter $dz$. 

**Fig. 3.7** Changes in patient dose due to increased beam filtration at constant contrast-to-noise ratio for different types of detail. Standard filtration: 1 mm Al + 0.1 mm Cu (= 4.5 mm Al quality equivalent); added filtration: 0.2 mm Cu (= 7 mm Al quality equivalent).
Dynamic range requirements for the detector system can thus be reduced. Simultaneously, beam-hardening effects are also less likely.

In order to provide attenuating properties that are almost tissue equivalent, beam shapers should be made from materials containing only elements with a low atomic number Z. However, this is not always the case in practice. Beam shapers preferentially affect the dose in the outer portions of an object, thereby reducing the peripheral CTDI values. But as the dose at the centre is mainly caused by scattered radiation from the periphery of the object, the central CTDI value is also somewhat reduced. The ratio of dose at the periphery to the dose at the centre therefore decreases, making the dose distribution inside an object more homogeneous and so improving the uniformity of noise in the image. Contrary to the flat filter, the beam shaper has a much greater impact on the dose properties of a scanner.

The beam shapers found in practice not only differ by the material from which they are made. They also differ by their shape, thus producing more or less compensation. A prominent example is the beam shaper of the Elscint CT Twin that was modified in 1998 to produce more compensation. In addition, different types of beam shapers can be selected on some scanners, depending on the nature and diameter of the object (e.g. for head and body scanning mode).

### 3.2.3 Beam Collimation

The beam collimation for defining the thickness of the slice to be imaged is made in the first instance close to the X-ray source (primary collimation). The shape of the dose profile is determined by the aperture of the collimator, its distance from the focal spot, and the size and shape (i.e. the intensity distribution) of the focal spot. Due to the narrow width of collimation, penumbral effects occur. These effects become more and more pronounced as collimation is further narrowed.

In addition, there is a secondary collimation close to the detector (‘post-patient collimation’) that primarily serves to remove scattered radiation. On some single-slice and dual-slice scanners this secondary collimation is further narrowed in order to improve the shape of the slice profile (‘restrictive post-patient collimation’, see fig. 3.8a,b). For multi-slice scanners with more than two detector rows, the primary collimation must necessarily be made wider than N times the selected slice collimation in order to avoid (or at least to reduce) penumbral effects in the outer portions of the detector array (fig. 3.8c). In both cases, the dose profile is wider than the slice profile or the nominal beam width, and the patient is exposed to a larger extent (‘overbeaming’), as becomes obvious from normalized CTDI values that increase with reduced beam width.

Overbeaming can be expressed by a single parameter, the ‘overbeaming parameter’ \( dz \), that is equal to the combined width of the portion of the dose profile that is not used for detection (fig. 3.8c). Overbeaming itself, i.e. the percentage increase in CTDI due to the unused portion of the dose profile, is then given by

\[
\Delta \text{CTDI}_{\text{rel}} = \frac{dz}{Nh_{\text{coll}}} \cdot 100 \quad (3.14)
\]

The overbeaming parameter \( dz \) typically amounts to 1 mm
for single- and dual-slice scanners that employ restrictive post-patient collimation, and to 3 mm for multi-slice scanners with \( N = 4 \) and more slices that are acquired simultaneously, but may vary depending on the type of scanner. For narrow beam width settings the increase in dose that results from overbeaming can be 100% and more.

In practice, overbeaming is no real issue for single- and dual-slice scanners, as the limited coverage restricts the use of narrow beam width to a few examinations with a short scan range (e.g. inner ear). With multi-slice scanners, however, overbeaming effects have to be taken seriously, as MSCT technology aims to provide improved resolution along the z-axis, which requires reduced slice collimation. Overbeaming, i.e. the increase in CTDI that results from beam width settings that are typical for each type of scanner is shown in fig. 3.9 for a number of scanners from different manufacturers. As indicated by the trend line, overbeaming is most pronounced with quad-slice scanners and is diminished with an increasing beam width \( N \cdot h_{\text{col}} \) provided by scanners with more slices (Nagel 2005).

### 3.2.4 Detector Array

In contrast to single-slice scanners, multi-slice scanners are equipped with a detector array that consists of more than a single row of detectors. Gas detectors or fourth-generation stationary detector rings are no longer compatible with multi-slice requirements. Consequently, only third-generation detector arcs with solid-state detectors have remained. In general, solid-state detectors are more dose-efficient than gas detectors (van der Haar et al. 1998), but require additional means to suppress scattered radiation (anti-scatter-grids) that inevitably cause a certain loss of primary radiation, too.

The single detectors in a multi-row, solid-state detector array are separated by narrow strips (‘septa’) which are not sensitive to radiation and therefore do not contribute to detector signal. Due to the large number of additional strips, these inactive zones result in minor or major geometrical losses, depending on the design of the detector array. In addition, further losses occur due to a decrease in sensitivity at the edges of each row that results from cutting the scintillator crystal. In contrast to a single-row detector array whose width can be larger than the maximum slice thickness (see fig. 3.10), the edges of the rows in a multi-row detector array are located inside the beam. Due to both these effects - separating strips and decreased sensitivity - the net efficiency of a solid-state detector array, which is typically 85% for single-slice scanners, is further decreased to typically 70%.

When 4-slice scanners were introduced in 1998, very different detector designs were used (fig. 3.11), with variations in the number of rows (between 8 and 34) and the smallest detector size (between 0.5 and 1.25 mm). The large number of rows (much larger than the number \( N \) of slices that can be acquired simultaneously) was necessary to enable the use of different slice collimations (between 4-0.5 mm and 4-8 mm). Slice collimations wider than the detector size are achieved by electronically combining several adjacent detector rows (e.g. 4-1.25 mm = 5 mm
Equipment-related Factors

Each detector design had its specific advantages and drawbacks: Toshiba’s hybrid arrangement offered the largest coverage (32 mm) and the acquisition of four sub-millimetre slices, but had the largest number of septa (1 per mm) and the smallest detector size (0.5 mm). The progressive design, commonly used by Philips and Siemens, had the smallest number of septa (0.35 per mm), but was restricted to two sub-millimetre slices only. GE’s matrix arrangement was a compromise (0.75 per mm) that, however, facilitated the next technology step towards eight simultaneously acquired slices with the same detector array.

All 16-slice scanners introduced in 2001 now made use of the same hybrid design, with 16 smaller central detectors, accompanied by a number of larger detectors at both sides (fig. 3.12). Apart from the number of detector rows (between 24 and 40) and array width (between 20 and 32 mm), there were differences in the size of the detectors (between 0.5 and 1.5 mm), and each manufacturer claimed his solution to be the best one. As in real life, there are a number of conflicting needs (spatial resolution, dose efficiency, coverage) that must be met, especially with respect to cardiac imaging where scan times below 20 s (one breathhold) are mandatory. Consequently, designs, which put emphasis to a single one of these criteria, only were definitely not the best compromise. Due to the increased number of septa (from 0.6 per mm (4-slice) to 1.1 per mm (16-slice) on average), the geometric efficiency of 16-slice detector arrays is somewhat lower.

In the latest generation of 64-slice scanners, matrix arrangements that allow for simultaneous acquisition of 64 sub-millimetre slices are employed by the majority of manufacturers (fig. 3.13). By electronically combining several adjacent rows, thicker slices can be acquired, too, but at a reduced number of slices (e.g. 32·1.25 mm, 16·2.5 mm etc.). Once again, the number of septa was increased (to 1.6 per mm on average), resulting in an additional loss in geometric efficiency.

The hybrid detector design exclusively used by Siemens for its Sensation 64 scanner is particular insofar as the number of simultaneous slices claimed by the manufacturer (64) is much larger than the number of rows (32·0.6 mm or 24·1.2 mm). The claim is based on a special acquisition mode that employs two alternating focal spot positions to simultaneously produce 64 data sets per rotation with...
50% overlap in order to achieve a somewhat improved spatial resolution in z-direction. With respect to all other important features (collimation, coverage, overbeaming effects etc.), however, this model behaves as a 32-slice scanner in submillimetre mode and a 24-slice scanner in all other modes at maximum. In addition, the thickness of the smallest slice that can be reconstructed (relevant for partial volume effects) is at least equal to the smallest slice collimation, i.e. 0.6 mm (Flohr et al. 2004), not lower.

### 3.2.5 Data Acquisition System

The data acquisition system (DAS) serves to collect the detector signals, to convert them into digital information and to transfer the data to the image reconstruction system. The number of DAS channels, not the number of detector rows, is the decisive parameter that limits the number \( N \) of independent slices that can be acquired simultaneously. Consequently, the term ‘MDCT (multi-detector-row CT)’ is somewhat misleading, as has recently become the term ‘MSCT (multi-slice CT)’, too. So ‘multi-channel CT (MCCT)’ would be the most unequivocal notation.

With the advent of 16-slice scanners at latest, the spatial requirements of an increased number of detector rows and the exorbitantly increased data rate no longer allowed to use traditional circuit boards. Instead, application-specific integrated circuits (ASIC’s) were developed, with significantly reduced dimensions (fig. 3.14) and drastically increased data transfer capabilities. As these AICS’s operate with reduced electronic noise, they are advantageous with respect to the dose efficiency of the detector assembly. This is demonstrated by fig. 3.15 where the dose that is necessary to obtain images of equal image noise at equal slice thickness was reduced by 25% with the introduction of this advanced DAS chip (Vlassenbroek 2004).

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**Fig. 3.14** The spatial requirements of an increased number of detector rows and the exorbitantly increased data rate necessitated the development of data acquisition systems with tiny application-specific integrated circuits (ASIC’s) that replaced the traditional circuit boards (courtesy: Philips Medical Systems).

![Fig. 3.14](image)

**Fig. 3.15** As a side effect, the lower electronic noise of ASIC-based data acquisition systems introduced with 16-slice scanners allows for a 25% dose reduction at constant noise (beam width for all scanners: \( N \cdot h_{vol} = 10 \text{ mm} \)).

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![Fig. 3.16](image)

**Fig. 3.16** The most common interpolation schemes for single-slice scanners are either 360° LI (a.) or 180° LI (b.). Both schemes employ two data points closest to the position \( z_0 \) of the reconstructed slice for each projection angle. Making use of the virtual (complementary) data (dashed lines), a shorter interpolation distance is achieved with 180° LI, resulting in a narrower slice profile.

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![Fig. 3.16](image)
3.2.6 Spiral Interpolation

Data acquisition in spiral scanning mode requires an additional interpolation step to obtain axial slices. The interpolation scheme of single-slice scanners employs two data points for each projection angle only, thus producing a bell-shaped slice profile. Depending on whether only true data (360° linear interpolation (LI), fig. 3.16a) or also virtual data (180° LI, fig. 3.16b) are used, the width of the slice profile is significantly broadened with increasing pitch (fig. 3.18a). The relative noise, however, remains independent from pitch and amounts to 83% (360° LI) and 117% (180° LI) compared to sequential scanning.

Most multi-slice scanners make use of a different interpolation scheme with more than two data points (‘z-filtering’, Taguchi and Aradate 1998). Depending on the slice thickness $h_{\text{rec}}$ to be reconstructed, interpolation is made using all data points that are located inside the pre-selected filter width $FW$ (= $h_{\text{rec}}$, fig. 3.17). Contrary to single-slice scanners, the width of the slice profile thus remains unaffected from changes in pitch settings of up to $p=2$ (fig. 3.18b). However, as the number of data points inside $FW$ is reduced, the noise will increase with pitch unless corrective actions are taken. This can be accomplished by adjusting the (electrical) tube current $I_{\text{el}}$ proportional to the change in pitch $p$. This adjustment is automatically made for all MSCT scanners from Elscint, Philips and Siemens, thereby using a different mAs notation ($Q_{\text{eff}}$) named ‘effective mAs’ or ‘mAs per slice’ that is different from the traditional electric mAs product $Q_{\text{el}}$:

$$Q_{\text{eff}} = \frac{I_{\text{el}} \cdot t_{\text{ms}}}{p} = \frac{Q_{\text{el}}}{p} \quad (3.15)$$

As a result, pitch has no longer any influence on slice profile width, image noise and average dose (CTDI vol) if $Q_{\text{eff}}$ is held constant. This does not hold for MSCT scanners manufactured by General Electric and Toshiba, which do not automatically correct the tube current for pitch and do not use effective mAs notation.

3.2.7 Adaptive Filtration

Adaptive filtration (AF) is a dedicated data processing technique for projections that are subject to strong atten-
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3.2.8 Overranging

‘Overranging’ is the increase in dose-length product due to the additional rotations at the beginning and at the end of a spiral scan required for data interpolation to reconstruct the first and the last slice of the imaged body region. With single-slice scanners, theory requires that an additional rotation is usually made in total (Kalender 2000). For multi-slice scanners, the situation is much less obvious, as will be seen from the results presented below.

Overranging effects can be expressed both in terms of the additional number $\Delta n$ of rotations and the increase $\Delta L$ in scan length. $\Delta L$ primarily depends on two factors: the beam width $N \cdot h_{\text{col}}$ and the pitch factor $p$. This can be fairly well described by a linear relationship (Nagel 2005):

$$\Delta L = (m_{\text{sp}} \cdot p + b_{\text{sp}}) \cdot N \cdot h_{\text{col}}$$

(3.16)

While single-slice scanners behave as expected from theory, the characteristics of typical MSCT scanners differ mark-

![Fig. 3.19 Projections suffering from excessive attenuation result in images with unisotropic noise patterns (left); images processed with adaptive filtration show a reduced and more homogeneous noise pattern (right).](image)

![Fig. 3.20 Adaptive filtration affects only those projections where the attenuation exceeds a pre-selected level.](image)

![Fig. 3.21 While single-slice scanner (SSCT) usually require only one additional rotation $\Delta n$ in spiral scanning mode, multi-slice scanners (MSCT) show a pronounced pitch dependence. Conversely, the normalized elongation of the scan range, $\Delta L/N \cdot h_{\text{col}}$, is almost constant for most MSCT scanners, but increases linearly with pitch for SSCT scanners.](image)
Equipment-related Factors

The number $\Delta n$ of additional rotations (fig. 3.21a) is strongly pitch dependent, while the normalized elongation of the scan range, $\Delta L/N \cdot h_{col}$, is almost independent from pitch (fig. 3.21b) and amounts to approximately 1.5, i.e. $\Delta L$ is typically 1.5 times the total beam width $N \cdot h_{col}$. For most single-slice scanners, the overranging parameters $m_{OR}$ and $b_{OR}$ are equal to 2 and −1, respectively. For the majority of MSCT scanners, typical values for $m_{OR}$ and $b_{OR}$ are 1 and 0.5, respectively.

The implications of overranging effects for the radiation exposure to the patient, i.e. the dose-length product DLP, not only depend on $\Delta L$, but also on the length $L_{net}$ of the imaged body region. The percentage increase in DLP is given by

$$\Delta DLP_{rel} = \frac{\Delta L}{L_{net}} \cdot 100 \quad (3.17)$$

Fig. 3.22 Overranging, i.e. the percentage increase in DLP, for single-slice ($N=1$), dual-slice ($N=2$), quad-slice ($N=4$), 6 to 8-slice ($N=6-8$), 16-slice ($N=16$) and 32 to 40-slice ($N=32-40$) scanners from different manufacturers (A to F) for a scan length $L$ of 20 cm and the slice collimations and pitch settings typically employed. The red trend line indicates that overranging becomes more pronounced with scanners that allow for a wider beam width $N \cdot h_{col}$.

The extent of overranging is shown in fig. 3.22 for a representative selection of single and multi-slice scanners from different manufacturers for typical scan parameter settings and a typical scan length of 20 cm. Overranging effects are normally almost negligible for single-slice and the majority of dual- and quad-slice scanners. Contrary to overbeaming, overranging becomes larger with an increasing number of slices acquired simultaneously due to the enlarged beam width. Even greater values might occur for beam widths larger than the typical ones assumed here and scan ranges being shorter than 20 cm.

3.2.9 Devices for Automatic Dose Control

Newer scanners are equipped with means that automatically adapt the mAs settings to the individual size and

Fig. 3.23 Automatic exposure control (AEC) accounts for the average attenuation of the patient’s body region that is to be scanned. For slim patients, mAs is reduced to a level that ensures constant image quality.

Fig. 3.24 Longitudinal dose modulation (LDM) is a refinement of AEC that adapts the mAs settings slice-by-slice or rotation by rotation. Those parts of the scan range with reduced attenuation will be less exposed.
shape of the patient. As this matter is discussed in detail in chapter 6, only a brief overview shall be given here.

Automatic dose control systems offer up to four different functionalities, which can be used either alone or in combination:

- Automatic exposure control (AEC, fig. 3.23) that accounts for the average attenuation of the patient’s body region that is to be scanned. Information on the patient’s attenuation properties is derived from the scan projection radiogram (SPR) usually recorded prior to the scan for planning purposes.
- Longitudinal dose modulation (LDM, fig. 3.24), which is a refinement of AEC by adapting the mAs settings locally, i.e. slice-by-slice or rotation by rotation.
- Angular dose modulation (ADM, fig. 3.25), another refinement of AEC that adapts the tube current to the varying attenuation at different projection angles. Information on the patient’s attenuation properties is either derived from two SPR or in real-time from the preceding rotation.
- Temporal dose modulation (TDM, fig. 3.26) that reduces the tube current in cardiac CT (or other ECG-gated CT examinations) during those phases of the cardiac cycle that are not suited for image reconstruction due to excessive object motion.

The common denominator of these functionalities is that the user no longer needs to select his parameter settings with respect to the ‘worst case’, i.e. obese patients, the part of the scan range with the highest attenuation (e.g. shoulder in chest exams), the projection with the highest attenuation (lateral) etc.. Consequently, a significant dose reduction from the application of these devices can be expected.

All major CT manufacturers now offer some or all of these functionalities with their latest scanners. A comprehensive report on the current status of automatic dose control systems has been published by ImPACT (2005). However, there are significant differences how these devices operate and perform. At present, some of these systems are not sufficiently user-friendly and make adjustments in a way that seems to be theoretically sound, but does not comply with other, more comprehensive aspects of image quality. Some of these shortcomings will be discussed in the following section.

### 3.2.10 Dose Display

Newer scanners must be equipped with a dose display. At present, only the display of CTDI\textsubscript{vol} is mandatory (IEC 2001). However, many scanners already show the DLP, too, either per scan series or both DLP per scan series and DLP per exam. An example with display of CTDI\textsubscript{vol} and DLP per scan series is shown in fig. 3.27.

With the dose display, dose is not saved per se, but feedback is provided that may help to achieve this goal, e.g.
by comparison of the displayed dose values with dose recommendations. In addition, changes in scan parameter settings and their implications for patient exposure are made immediately obvious. Thus the dose display can be used for purposes of dose optimisation. Finally, CTDI$_{vol}$ can be used as a fair estimate for the dose to organs that are entirely located in the scan range.

The interpretation of the dose values displayed at the scanner’s console needs special attention in the following situations:

- Many dose recommendations are given in terms of weighted CTDI (CTDI$_w$); in order to allow for comparisons, the pitch correction involved in CTDI$_{vol}$ must be reverted by multiplying CTDI$_{vol}$ with the pitch factor.
- Up to now, the dose values for examinations carried out in body scanning mode are always based on body-CTDI regardless of patient size. In paediatric CT examinations, the displayed figures should be multiplied by 2 for children and by 3 for infants in order to give a realistic estimate of patient dose.

**Fig. 3.27** Scan protocol window of a Philips Mx8000 IDT scanner with dose display (CTDI$_{vol}$ and DLP per scan series) at the bottom.
### 3.3 Application-related Factors

Although the scanner design is of some importance, surveys on CT practice have regularly shown that the way how the scanner is used has the largest impact on the doses applied in a CT examination. The application-related factors on which patient exposure depends are subdivided in

- **scan parameters**, i.e. those factors that directly determine the local dose level ($\text{CTDI}_{\text{vol}}$) and that are often pre-installed or recommended by the manufacturer (e.g. in application guides);
- **examination parameter**, i.e. those factors that – in combination with $\text{CTDI}_{\text{vol}}$ - determine the integral exposure (i.e. DLP) and that depend on the preferences of the user;
- **reconstruction and viewing parameters**, which implicitly influence the dose settings.

First, however, the principal interdependences between dose settings and image quality shall be outlined

#### 3.3.1 Brooks’ Formula

As in conventional projection radiography, aspects of dose and image quality are linked. For CT, Brooks and DiChiro (1976) have formulated the correlation between these two opposed quantities:

$$D \propto \frac{B}{\sigma^2 \cdot a^2 \cdot b \cdot h} \quad \text{with } B = \exp^{-\mu d} \quad (3.18)$$

where

- $D$ = patient dose
- $B$ = attenuation factor of the object
- $\mu$ = mean attenuation coefficient of the object
- $d$ = diameter of the object
- $\sigma$ = standard deviation of CT numbers (noise)
- $a$ = sample increment
- $b$ = sample width
- $h$ = slice thickness

This fundamental equation - commonly known as the ‘Brooks’ formula’ - describes what happens with respect to patient dose if one of the parameters is changed while image noise remains constant:

- dose must be doubled if slice thickness is cut by half;
- dose must be doubled if object diameter increases by 4 cm;
- an eightfold increase in dose is required if spatial resolution is doubled (by cutting sample width and sample increment by half).

In this context, the term ‘dose’ is applicable to each of the dose quantities that are appropriate for CT. Dose and noise are inversely related to each other in such a way that a fourfold increase in dose is required if noise is to be cut by half.

It should be noted, however, that the Brooks’ formula is incomplete in that image quality is only considered in terms of quantum noise and spatial resolution. Other important influences, such as contrast, electronic noise or artefacts, are not taken into account and will therefore modify optimization strategies under particular circumstances.

#### 3.3.2 Scan Parameters

##### 3.3.2.1 TUBE CURRENT-TIME PRODUCT ($Q$)

As in conventional radiology, a linear relationship exists between the tube current-time product and dose; i.e. all dose quantities will change by the same amount as the applied mAs. The mAs product $Q$ for a single sequential scan is obtained by multiplying the tube current $I$ and exposure time $t$; in spiral scanning mode, $Q$ is the product of the tube current $I$ and rotation time $t_{\text{rot}}$. This should not be mixed up with the total mAs product of the scan which is the product of tube current $I$ and (total) scan time $T$.

The consequences on image quality resulting from variations in the tube current-time product are relatively simple to understand. The only aspect of image quality so affected is image noise, which is - as indicated in equation (3.18) - inversely proportional to the square root of dose (i.e. mAs).

The tube current-time product is often used as a surrogate for the patient dose (i.e. CTDI). However, this is highly misleading, as the normalized CTDI values and thus the dose that results for the same mAs setting can vary by up to a factor 6 between different scanners. So it makes absolutely no sense to communicate dose information or recommendations on the basis of mAs. Instead, only $\text{CTDI}_{\text{vol}}$ (and DLP) should be used for this purpose.

With the advent of multi-slice scanners, additional confusion arose due to the introduction of a different, pitch-corrected mAs notation (‘effective mAs’ or ‘mAs per slice’, see equation (3.15)) by Elscint, Philips and Siemens. As most multi-slice scanners make use of a multi-point spiral interpolation scheme as outlined in section 3.2.6, effective mAs is the most appropriate notation for MSCT. Nevertheless, General Electric and Toshiba still prefer the traditional electrical mAs notation which fur-
ther makes it difficult to compare mAs settings from different scanners. This particularly holds for cardiac CT where very low pitch settings are used.

**Recommendation**

The settings for the tube current-time product should be adapted to the characteristics of the scanner, the size of the patient (see section 3.3.2.5), and the dose requirements of each type of examination. Examinations with high inherent contrast, such as for chest or skeleton, that are characterised by viewing with wide window settings, can regularly be conducted at significantly reduced mAs settings.

### 3.3.2.2 TUBE POTENTIAL (U)

When the tube potential is increased, both the tube output and the penetrating power of the beam are improved, while image contrast is adversely affected. In conventional projection radiography, increased tube potentials are applied in order to ensure short exposure times for obese patients, to equalize large differences in object transmission (e.g. during chest examinations) or to reduce patient dose. In the latter case, automatic exposure control (AEC) guarantees that the improved penetrating power of the beam is exclusively for the benefit of the patient.

In CT, increased tube voltages are used preferentially for improvements in tube loading and image quality. Contrary to the case for mAs, the consequences of variations in kV cannot easily be assessed. The relationship between dose and tube potential U is not linear, but rather of an exponential nature which varies according to the specific circumstances. The intensity of the radiation beam at the detector array, for example, varies with U to the power of 3.5. If the tube potential is increased e.g. from 120 to 140 kV, the electrical signal obtained from the detectors therefore changes by a factor 1.7 (fig. 3.28).

The decrease in primary contrast which normally results from this action is largely over-compensated by the associated decrease in noise, i.e. the higher the tube potential, the better the contrast-to-noise ratio CNR (except for the application of iodine as contrast agent). The only reason why this analysis generally holds true is the absence of any kind of AEC in the majority of scanners which might prevent unnecessary increases in the detector signal. This clearly demonstrates that dose is not reduced by applying higher kV settings, but merely increased as long as mAs settings are not changed: weighted CTDI and effective dose increase with U to the power of 2.5 (fig. 3.28), which means that both are increased by approximately 50% if kV settings are changed from 120 to 140 kV.

Therefore the question is justified whether and when it might be reasonable to deviate from the 120 kV setting usually applied. As can be seen from fig. 3.29, this depends on the attenuation characteristics of the detail that is diagnostically relevant. The figures are given in terms of contrast-to-noise ratio squared (CNR²) at constant patient dose; this notation allows to directly convert the percentage differences into dose differences. For soft tissue contrast (e.g. differences in tissue density), higher tube potentials perform slightly better than lower ones, but the differences are quite small. The opposite holds true for bone contrast (i.e. bone vs. tissue). For iodine contrast, however, there is a strong dependence on tube potential that is much in favour of lower kV settings. So 80 instead of 120 kV would allow to reduce the patient dose by almost a factor of two without sacrificing image quality.
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**Recommendation**

Tube potentials other than 120 kV should be considered only in case of

- obese patients where mAs cannot further be increased: use higher kV settings
- slim patients and paediatric CT where mAs cannot further be reduced: use lower kV settings
- CT angiography with iodine: use lower kV settings.

Variations in tube potential should not be considered for pure dose reduction purposes except for CT angiography. Due to the complexity involved, adaptation of mAs settings should not be left to automatic exposure control systems, as these do not account for changes in contrast. Dose settings in CT angiography should not be higher than in unenhanced scans of the same body section and should be lowered if performed at reduced kV settings.

### 3.3.2.3 SLICE COLLIMATION ($h_{col}$) AND SLICE THICKNESS ($h_{rec}$)

With single-slice CT, the slice collimation $h_{col}$ used for data acquisition and the reconstructed slice thickness $h_{rec}$ used for viewing purposes were identical (except for slice profile broadening in spiral scans with increased pitch as discussed in section 3.2.6). So there was no need to distinguish between both of them. With multi-slice CT, the slice collimation (e.g. 0.75 mm) and the reconstructed slice thickness (e.g. 5 mm) are usually different. Frequently, the selection of the reconstructed slice thickness is made with respect to multiplanar reformatting (MPR) purposes (e.g. 1 mm), thus creating a so-called ‘secondary raw data set’, i.e. a stack of thin slices from which MPR slabs with larger thickness (e.g. 5 mm) can be made for viewing purposes.

The ability to acquire longer body sections with thin slices in order to achieve an almost isotropic spatial resolution is the most important achievement of multi-slice technology. As reduced slice thickness is associated with increased image noise, this may have a significant impact on patient dose as expressed by the Brooks’ formula (equation 3.18). Therefore it is worth while to treat this matter in a somewhat more detailed fashion.

A narrow slice collimation is a precondition for a narrow slice thickness, but its impact on patient dose is restricted to aspects of overbeaming and overranging only. As these show opposed dependences on beam width, as outlined in sections 3.2.3 and 3.2.8, the question arises for the optimized beam width settings. As demonstrated for a typical MSCT scanner in fig. 3.30, beam width settings greater than 10 mm perform almost equally well (a.) except for short scan ranges (spine, paediatrics) where a beam width

---

**Fig. 3.30** Increased dose-length product due to overbeaming (OB) and overranging (OR) effects for a typical MSCT scanner. For average to long scan ranges (L = 20 cm and more, a.), all beam width settings above 10 mm perform almost equally well. For short scan ranges (L = 10 cm as in paediatric and spine exams, b.), beam width settings between 10 and 20 mm should be preferred.

**Fig. 3.31** Relative image quality in dependence of the slice thickness $h_{rec}$. Improvements in image quality (better detail contrast due to reduced partial volume effects) outweigh the deteriorations caused by increased noise. As a result there is a net gain in contrast-to-noise ratio (CNR) at reduced slice thickness without any increase in dose.

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The relative image quality in dependence of the slice thickness $h_{rec}$. Improvements in image quality (better detail contrast due to reduced partial volume effects) outweigh the deteriorations caused by increased noise. As a result there is a net gain in contrast-to-noise ratio (CNR) at reduced slice thickness without any increase in dose.
of between 10 and 20 mm is more appropriate (b.). Beam width settings below 10 mm should be avoided due to increased overbeaming effects unless there are other important aspects that justify to override this recommendation.

The decisive determinant with respect to image noise and its implications for patient dose, however, is the slice thickness that is finally used for viewing purposes. The relationship between slice thickness, noise and dose expressed in the Brooks’ formula tends to correct any reduction in slice thickness by a corresponding increase in dose to ensure a constant image noise, and some automatic exposure control systems exactly do so. However, any variation in slice thickness also affects image contrast due to a modification in partial volume effect, which is not taken into account by the Brooks’ formula. As shown in fig. 3.31, image noise and image contrast of small details will react in a different fashion on reduction of the slice thickness: while image quality in terms of noise is impaired

Fig. 3.32 MSCT examination of the liver performed on a MSCT scanner (Siemens Somatom Volume Zoom) at 120 kV, 4-2.5 mm slice collimation and 125 mAs (CTDI vol = 11 mGy). From the same raw data set, slices of different thickness (3 mm (a.), 5 mm (b.), 7 mm (c.), and 10 mm (d.)) were reconstructed at the same central position z. Despite the increased noise pertaining for thinner slices, the visibility of small lesions improves remarkably owing to reduced partial volume effects. This is clearly demonstrated by a lesion approximately 3 mm in size (arrow) (courtesy Dr. Wedegaertner, University Hospital Eppendorf, Hamburg, Germany).
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proportional to the square root of the change in slice thickness only, the contrast is improved proportional to the slice thickness. As a result, there is a net gain in image quality in terms of contrast-to-noise ratio CNR without any increase in dose whenever partial volume effect is of importance.

This is clearly demonstrated by the clinical example given in fig. 3.32, where the visibility of a liver lesion (approximately 3 mm in size) diminishes continually with increasing slice thickness – despite reduced image noise. In addition, a detailed analysis of the results of the German survey on CT practice in 1999 (Galanski et al. 2001) has revealed that slice thickness has only minor or no influence on clinical dose settings. This is shown in fig. 3.33 for liver examinations with slice thicknesses of between 3 and 10 mm that were used in practice. Therefore it is essential to understand that the selection of a narrow slice collimation is only a means to an end: to enable MPR images without or with reduced step artefacts, and, if necessary, to overcome partial volume effects.

**Recommendation:**
The slice collimation should be selected as small as compatible with aspects of overbeaming/overranging, total scan time and tube power. Viewing should preferably be made with thicker slabs (e.g. 3 to 8 mm), thereby reducing image noise and other artefacts. Thinner slabs should only be used if partial volume effect is of importance. This should preferentially be done in conjunction with workstations that allow to change the slab thickness in real-time. Except for very narrow slices there should be no need for any increase in dose settings on reduction of slice thickness.

3.3.2.4 PITCH (p)

With SSCT scanners, scanning at increased pitch settings primarily serves to increase the speed of data acquisition. As a side effect, patient dose is reduced accordingly, at the expense of impaired slice profile width, i.e. z-resolution, however. As already outlined in section 3.2.6, MSCT scanners make use of a spiral interpolation scheme that is different from SSCT. Thus the slice profile width remains unaffected from changes in pitch settings. Instead, image noise changes with pitch (fig. 3.34a) unless the tube current is adapted accordingly.

Scanners that make use of the effective mAs (mAs per slice) concept not only keep slice profile width, but also image noise constant when pitch changes (fig. 3.34a). To achieve this goal, the electrical mAs product supplied to the x-ray tube automatically changes linearly with pitch (fig. 3.34b). As a consequence, patient dose (CTDIvol) is no longer reduced at increased pitch settings in contrast to SSCT scanners. On the other hand, dose will also not increase at reduced pitch settings. MSCT scanners without automatic adaptation of mAs will still save dose at increased pitch setting, but this will happen at impaired

![Fig. 3.34 For MSCT systems that employ multi-point spiral data interpolation (z-filtering), image noise changes with pitch unless effective mAs is held constant (a.). This implies that the electrical mAs product supplied to the tube changes with pitch (b.). Contrary to SSCT, changes in pitch settings therefore no longer have any influence on patient dose in terms of CTDIvol.](image-url)
image quality (more noise) as long as mAs is not adapted manually.

Frequently, image quality in terms of artefacts depends on pitch settings. In general, spiral artefacts are reduced at lower pitch settings. For similar reasons, some scanners allow the setting of a limited number of ‘preferred’ pitches only. Reduced pitch settings can also be applied to enhance the effective tube power, however, at the expense of reduced scanning speed.

Recommendation:
Pitch settings with MSCT scanners should be made exclusively with respect to scan speed, spiral artefacts and tube power. Dose considerations no longer play a role if scanners that employ effective mAs are used or if (electrical) mAs is adapted to pitch to achieve constant image noise.

3.3.2.5 OBJECT DIAMETER (d)

Patient size, although not a parameter to be selected at the scanners’s console, represents an important influencing parameter that needs to be considered in this context. Considerable reductions in mAs settings are appropriate whenever slim patients, and particularly children, are examined. In order to avoid unnecessary over-exposure, the mAs must be intentionally adapted by the operator unless AEC-like devices are available. Due to the decreased attenuation for the smaller object, image quality will not be impaired if mAs is selected appropriately. This means that the image quality will be at least as good as for patients of normal size, although the dose has been reduced.

The two questions to be solved in this context are:

• To which degree shall mAs settings be adapted in dependence of the object diameter d?
• Which diameter is typical for a standard patient to whom the standard protocol settings refer to?

From theoretical considerations (half-value thickness HVL for CT beam qualities), mAs should be altered by a factor 2 for each change in patient diameter of 4 cm tissue-equivalent thickness. However, dedicated studies (e.g. Wilting et al. 2001) have shown that this algorithm doesn’t work well in practice: Although objective (i.e. measured) noise was almost constant for patient diameters of between 24 and 36 cm, it was found that the subjective (i.e. perceived) image quality continually decreased with the patient diameter and vice versa. This is most likely due to the circumstance that adipose patients have more fatty tissue around their organs. Thus the inherent contrast is better, and more noise can be tolerated. The opposite holds true with slim patients.

Consequently, a more gentle adaptation of mAs with patient diameter (factor 2 in mAs per 8 cm change in patient diameter) will better comply with clinical needs. Among the automatic exposure control systems currently in use, those from Philips and Siemens already make use of this modified algorithm that ensures a constant ‘adequate’ image quality, while those implemented by General Electric and Toshiba simply attempt to ensure a constant noise level. As already outlined in sections 3.3.2.2 for tube potential and 3.3.2.3 for slice thickness, strategies for automatic dose control that do not account for image contrast will fall short with respect to clinical needs. Similar consideration apply to the longitudinal dose modulation functionality: in examinations comprising several consecutive body sections with differing attenuation properties (e.g. in tumor staging of chest, abdomen and pelvis in a single spiral acquisition), mAs adjustment is often made in a way that ensures constant image noise, thus producing the

Fig. 3.35 Relationship between patient weight and lateral diameter according to a detailed analysis of patient data from a big children’s hospital (a.) and relative mAs settings in dependence of patient weight as recommended by three representative authors (b.). As indicated by the dashed lines, mAs adaptation by a factor of 2 per 8 cm change in patient diameter is almost perfectly met by Rogalla’s recommendation.
highest settings in the pelvis region. However, inherent contrast in the pelvis region is much better than in the upper abdomen; consequently, reduced mAs settings would be more appropriate, as recommended in ICRP publication 88 (ICRP 2001).

Although not specified explicitly, standard protocol settings implemented by the manufacturers are usually tailored to satisfy the vast majority of clinical situations except for obese patients where higher mAs or kV settings must be applied. So there is good reason to refer these standard settings to patients of about 80 to 85 kg body weight, which also is the average weight of European males. This corresponds to a lateral diameter of 33 cm according to a detailed analysis of patient data from a large children’s hospital in Germany (Schneider 2003, fig. 3.35a.). The following formula can be used to convert from lateral patient diameter \( \text{d}_{\text{lat}} \) (in cm) to patient weight \( \text{m} \) (in kg) and vice versa:

\[
\text{d}_{\text{lat}} = 6.5 + 3 \cdot \sqrt[3]{\text{m}} \quad (3.19)
\]

In current literature, numerous differing recommendations can be found on how to reduce mAs settings with patient weight or diameter. In fig. 3.35b, three examples are shown which are representative for a weak (Donelly et al. 2001), moderate (Rogalla 2004) or strong (Huda et al. 2000) adaptation of mAs to patient weight. As indicated by the dashed lines, mAs adaptation by a factor of 2 per 8 cm change in patient diameter is almost perfectly met by Rogalla’s recommendation which follows a very simple relationship:

\[
\text{Relative mAs} \propto \text{body weight} + 5 \text{ kg} \quad (3.20)
\]

A similar relationship has been proposed by an other research group (Honnef et al 2004). This formula can be used to create a set of standard protocols for different weight classes (e.g. 0-10 kg, 11-20 kg, 21-40 kg, 41-60 kg, 61-80 kg etc.) which can easily be applied in daily practice.

**Recommendation:**

mAs settings should be adapted to patient size in a more gentle way (factor 2 per 8 cm change in diameter) than predicted by theoretical considerations that only account for image noise. In addition, body regions with better inherent contrast should be scanned at reduced mAs settings. Preferentially, AEC systems that rather measure than adapt mAs to patient weight. As indicated by the dashed lines, mAs adaptation by a factor of 2 per 8 cm change in patient diameter is almost perfectly met by Rogalla’s recommendation which follows a very simple relationship:

\[
\text{Relative mAs} \propto \text{body weight} + 5 \text{ kg} \quad (3.20)
\]

A similar relationship has been proposed by a research group (Donnelly et al. 2001). This formula can be used to create a set of standard protocols for different weight classes (e.g. 0-10 kg, 11-20 kg, 21-40 kg, 41-60 kg, 61-80 kg etc.) which can easily be applied in daily practice.

3.3.3 Examination Parameters

3.3.3.1 SCAN LENGTH (L)

As already pointed out in section 3.1, the local dose, i.e. CTDI, is almost independent of the length of the scanned body section. The same does not hold, however, for the integral dose quantities, i.e. dose-length product and effective dose. Both increase in proportion to the length of the body section. Therefore, limiting the scan length according to the clinical needs is essential.

On most scanners, the scan length \( L \) is usually not indicated explicitly. Instead, the positions of the first and the last slice are stated only; the same holds for the information that is documented on the images or in the DICOM data file. The net scan length \( L_{\text{net}} \), i.e. the length of the imaged body section, is calculated by

\[
L_{\text{net}} = \left| pos. \text{ first sl} - pos. \text{ last sl} \right| + h_{\text{rec}} \quad (3.21)
\]

while the gross scan length \( L_{\text{gross}} \), i.e. the length of the irradiated body section, is

\[
L_{\text{gross}} = L_{\text{net}} + \Delta L \quad (3.22)
\]

where \( \Delta L \) is the increase in scan length due to overranging as described in equation (3.16). As a rule of thumb that holds for the majority of MSCT scanners, the actual scan range, overranging included, is extended at each side of the planned scan range by approximately 0.75·N·h rec + 0.5 h rec. This amounts to roughly 2 cm for a 16-slice scanner with 20 mm beam width and 5 mm slice thickness.

**Recommendation:**

For each patient, the scan length should be selected individually, based on the scan projection radiograph that is generally made prior to scanning for the purposes of localization, and should be kept as short as necessary. Moreover, a reduction in the scan range should be considered in multi-phase examinations and follow-up studies. Whenever feasible, critical organs like the eye lenses or the male gonads should be excluded from the scan range. This may be difficult for MSCT scanners that allow for large beam width settings due to increased overranging effects.

3.3.3.2 NUMBER OF SCAN SERIES (Ser)

In CT terminology, a scan series is usually referred to as a series of consecutive sequential scans or one complete spiral scan. With the limited tube power available for many SSCT scanners, CT examinations of long body sections (e.g. tumor staging of the entire trunk) had to be divided into several consecutive subsections. If the same protocol settings are applied to each series, the local dose will always be the same, while the integral dose is the sum of
the DLP or effective dose values of each series. So it would
not make a difference if the body section is scanned as a
whole or in several shorter subsections except for over-
ranging effects that will increase proportional to the num-
ber of subsections. On the other hand, mAs settings can
be adapted to the particular needs of each subsection, e.g.
lower settings for the chest, higher settings for the upper
abdomen and reduced settings for the pelvis, as indicated
in section 3.3.2.5.

If the same body section (or parts of it) is scanned more
than once, this is usually denoted as ‘multi-phasic’. How-
ever, this not only applies to examinations with administra-
tion of contrast agents, but also to examinations where
the same body section is scanned with different orientation
(like in facial bone exams) or with different slice collima-
tion settings (e.g. chest standard plus high resolution). Al-
though more than one scan is made at the same position,
the length of each single scan of a multi-phasic exam must
not necessarily be the same. While it is meaningful to sum
up the integral doses (DLP, effective dose) of each phase,
it is not for the local doses (i.e. CTDIvol). Nevertheless,
multi-phasic exams result in an increase in integral radi-
ation exposure that is roughly proportional to the number
of phases.

Recommendation:
The number of scan series (phases) should be kept as low
as necessary. This particularly hold for liver examinations
where studies with up to six different phases are sometimes
recommended in literature.

3.3.3.3 NUMBER OF ROTATIONS IN DYNAMIC CT
STUDIES (n)

In dynamic CT studies, e.g. in CT fluoroscopy or in
perfusion studies, a multiple number of scans is made at
the same position. Therefore it is meaningful to sum up
the local doses, too. For this particular situation, the main
issue is the avoidance of deterministic radiation effects.
Local doses can be quite high if the scans are made with
the standard dose settings used for that body region. In-
tegral doses are normally comparable to the values encoun-
tered in standard examinations of the same region. But
with the advent of wider detector arrays, which may be-
come even larger in future, integral dose will also signifi-
cantly be increased.

The doses applied in dynamic CT studies depend on two
factors: the dose, i.e. the CTDIvol, per rotation, and the
number n of rotations. As perfusion studies are regularly
made with administration of contrast agents, the benefits
of reduced kV settings as described in section 3.3.2.2
should be used to reduce the dose settings. The number n
of rotations can be kept low by limiting the total length of
the study, by reducing the image acquisition rate or by
intermitting the procedure (in CT fluoroscopy) whenever
possible.

Recommendations:
Dynamic CT studies should be made with the lowest dose
settings, the most narrow beam width, the shortest length
and the smallest image rate that is compatible with the
clinical needs of the examination.

3.3.4 Reconstruction and Viewing Parameters

3.3.4.1 FILTER KERNEL (FK)

CT images are reconstructed from sets of attenuation
measurements using dedicated mathematical procedures
(algorithms) which are known as ‘reconstruction filters’
or ‘filter kernels’. These algorithms are characterized as
having quite different properties with regard to image
quality: with highly-resolving filter kernels, spatial resolu-
tion is improved, but noise is increased. The opposite hap-
pens with smoothing kernels which reduce noise at the
expense of spatial resolution.

The properties of reconstruction filters are not subject to
standardization. Therefore, kernels of equal or similar
designation may vary considerably from one brand of
scanner to the next. Equally, reconstruction filters used
for head or body scans carrying the same name are by no
means identical. Labels such as ‘smooth’ or ‘sharp’ can
only be used as coarse indicators of the balance between
spatial resolution and image noise.

The compromise between spatial resolution and contrast
resolution for a particular clinical indication must be found
by appropriate selection of the reconstruction filter. The
better the spatial resolution, the higher the noise, as indi-

![Fig. 3.36 Typical noise characteristics of different filter kernels. Relative figures are given in terms of noise squared, so the percentage differences can directly be translated into dose differences that would be necessary for constant image noise.](image-url)
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cated in fig. 3.36. Image noise, however, strongly affects contrast resolution. Due to the relationship between dose and noise given by the Brooks’ formula (equation 3.18), the decision to use a particular filter kernel may directly affect the amount of dose required.

There are two practical sets of circumstances where dose can be saved by proper selection of the reconstruction filter. The first is where spatial resolution is more than sufficient for a given clinical indication. Contrary to the manufacturer’s recommendations, a smoother filter kernel can therefore be selected. The improvements resulting from this choice can then be used to reduce dose instead of noise, as indicated in fig. 3.37. The second is where the contrast-to-noise ratio for high contrast structures (e.g. lungs, skeleton) is more than sufficient, even though a highly-resolving filter kernel was used. In this case, increased noise can be tolerated, even if dose is reduced. So it turns out once again that the Brooks’ formula is somewhat misleading as it doesn’t account for image contrast. Nevertheless, one particular automatic exposure control system also tempts to compensate for changes in noise that result from the selection of the filter kernel.

**Recommendation:**
The selection of the filter kernel should be made with respect to the inherent contrast and as smooth as compatible with the clinical needs, thereby reducing the dose to

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**Fig. 3.37 Comparison of images that were scanned and reconstructed with different mAs and filter kernel settings but result in similar image noise.**

**Fig. 3.38 Comparison of two images with different settings of the window width W. Wider window settings result in smoother images, thus allowing for reduced dose settings, provided that the inherent contrast is sufficiently high.**
the noise level that is appropriate. High resolution kernels should only be used for high contrast structures without adaptation of mAs settings.

3.3.4.2 WINDOW WIDTH (W)

The window width is often not regarded as a relevant factor for influencing dose, since it is assumed that the width of the window is a parameter only related to image presentation. However, the visual perception of image noise strongly depends on the choice of window width setting. By using a wide window setting, noise perception can be reduced, as shown in fig. 3.38. The reduction is inversely correlated to window width (Prokop98). However, image contrast is also decreased, of course, because the number of grey scale values is simultaneously reduced.

Therefore a prerequisite for dose reduction by the use of wider window settings is sufficiency of contrast-to-noise ratio. Due to the non-linear relationship between dose and noise, even a relatively small increase in window width is profitable: if a setting of 350 HU is used instead of 300 HU, dose can be reduced by 26% while noise perception remains the same. It is therefore worthwhile finding out whether wider settings than those recommended by the manufacturer might also be appropriate. This holds particularly for high contrast structures; by doubling the window width, the dose can be cut to one quarter.

**Recommendation:**

The window width should be selected as wide as tolerable. With high contrast structures the improvement in noise thus achieved should be used to reduce the dose settings.

### 3.4 References

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Rogalla P (2004) Personal communication

Schneider K (2004) Personal communication


